

On Finding Better Wavelet Basis for Bearing Fault Detection

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Abstract: This paper considers the comparison of the Meyer and Morlet wavelet for bearing fault diagnosis. We created a wavelet based upon a transient vibration signal model established for signals generated in deep-groove ball bearings with pitting (spalling) formulation on their inner race. The wavelet creation used the sub-optimal algorithm devised by Chapa and Rao that matches a Meyer wavelet to a band limited signal in two steps. We tested the applicability of the matched wavelet for identifying this kind of bearing failure. The Morlet wavelet was used as a benchmark for evaluating the performance of the matched wavelet since many publications show its successful application. It was shown that for analysing exponentially or near-exponentially damped vibration responses like the vibration produced by spalling on the inner race of a deep-groove ball bearing, the Morlet wavelet is a reasonable choice and gives better results than the Meyer wavelet.

Keywords: Wavelet analysis; bearing vibration analysis; wavelet matching; condition monitoring

1 Introduction

It is known that the *Discrete Fourier Transform* (DFT) is most suitable for testing finite-energy, periodic, time-discrete quantities. The reason why the DFT is still used effectively for the vibration analysis of bearings is that most of the complex

vibrations show periodicity in time. This periodicity is closely related to the geometry and rotational speed of the bearing. Thus, the vibration components can be determined with reasonable accuracy.

Wavelet Analysis is a relatively new tool that has been successfully applied in many areas of science. In recent years, several researchers have proposed the use of wavelet transform to test bearing vibration signals where the FFT was ineffective [1-3]. Some scientists recommend the use of existing wavelet basis functions, while others create new wavelet bases. The proposed methods also differ in applied analysis techniques as well. These are *Continuous Wavelet Transform (CWT)*, *Discrete Wavelet Transform (DWT)*, *Wavelet-packet analysis*, *Matching Pursuit*, etc.

Junsheng et al. [4] used an impulse response wavelet to analyse faults in a roller bearing with CWT. Their wavelet is simply an exponentially damped sinusoid. Jiang et al. [5] proposed a hybrid method that combines the Morlet wavelet filter and sparse code shrinkage. Kankar et al. [6] compared three machine learning techniques for bearing fault diagnosis. These methods were the support vector machine (SVM), the artificial neural network (ANN) and self-organizing maps (SOM). For feature extraction they used a Meyer and Morlet wavelet. They found that the Meyer wavelet performs better with SVM classifier. Sheen [7] effectively applied the Morlet wavelet in the envelope detection for the vibration signal and found it also useful in the defect diagnosis of bearing vibrations. The application of the complex Morlet wavelet with SVM classifier is suggested in [8] for fault diagnosis of ball bearings having localized defects on various bearing components. Liu et al. [9] suggested an automatic feature extraction algorithm for bearing fault diagnosis using a correlation filter-based matching pursuit.

In wavelet analysis the choice of a wavelet is crucial from an analysis point of view. The analysing wavelet is usually independent of the signal investigated. Since the wavelet transformation and its derived energy distributions use convolution, one can obtain the highest output from these transformations when the signal and the wavelet are similar.

Over the past decade many publications [10-13] have considered creation of a matching wavelet to a given signal. *Tewfik* et al. [10] worked out a design method that matches a wavelet to the time domain form of a signal. *Chapa* and *Rao* [11] developed an algorithm that searches for a matching wavelet in frequency domain. Their method is capable of designing Meyer wavelets that approximate the wavelet amplitude and phase spectra separately. The cost function is the minima of the *Mean Squared Error (MSE)*, calculated from the amplitude spectra and group delay of the signal and the wavelet.

Our aim was to decide whether the Meyer or the Morlet wavelet is the better choice for analysing bearing failure.

We used *Chapa* and *Rao's* algorithm to create a matching wavelet to the signal model of transient vibration generated by an artificially created fault on the inner race of a deep groove ball bearing. To evaluate its performance, we used the vibration data of a pitted single row deep-groove ball bearing of type 6209. This bearing was earlier subjected to an endurance test. We calculated the scalogram of the vibration data using the newly created wavelet and the Morlet wavelet as well. Then we compared the results and found the better representation with Morlet wavelet.

2 Methods

2.1 Wavelets, Continuous Wavelet Transform and Scalogram

Wavelets [14] are functions constructed by translating and dilating a basic function called a mother wavelet Ψ (see Eq. (1)). The parameters a and b are called *scale* (dilation) and *translation parameters*, respectively. The wavelet is a normalised $\|\psi\| = 1$ function.

$$\psi_{b,a}(t) = \frac{1}{\sqrt{a}} \psi\left(\frac{t-b}{a}\right), \quad a > 0 \quad (1)$$

For $\Psi(t)$ to be a wavelet function and to recover $f(t)$ from its CWT, $\Psi(t)$ should satisfy some conditions. If $\Psi(t)$ has a zero average, i.e.:

$$\hat{\psi}(0) = \int_{-\infty}^{+\infty} \psi(t) dt = 0 \quad (2)$$

and satisfies the admissibility condition:

$$C_{\psi} = \int_{-\infty}^{+\infty} \frac{|\hat{\psi}(\omega)|^2}{|\omega|} d\omega < \infty, \quad (3)$$

where C_{ψ} is a constant that depends on the choice of wavelet, then there exists a *Continuous Wavelet Transform* (CWT) - *Inverse Wavelet Transform* (IWT) analysis-synthesis pair [15].

The CWT of a function $f(t)$ is defined as

$$CWT \{f(b, a)\} = \langle f, \psi_{b,a} \rangle = \int_{-\infty}^{+\infty} f(t) \cdot \frac{1}{\sqrt{a}} \psi^*\left(\frac{t-b}{a}\right) dt. \quad (4)$$

The benefit of CWT is that by changing the scale parameter, the duration and bandwidth of wavelet are both changed, providing better time or frequency resolution, but its shape still remains the same. The scale parameter can be linear or dyadic. The CWT uses short windows at high frequencies and long windows at low frequencies.

The *scalogram* [14], defined as the squared magnitude of CWT (Eq. (5)), always has non-negative, real-valued time-frequency (scale) distribution. This transformation conserves energy. Its resolution in the time-frequency plane depends on the scale parameter.

$$SC\{f(a, b)\} = |CWT\{f(a, b)\}|^2 = \left| \int_{-\infty}^{+\infty} f(t) \cdot \frac{1}{\sqrt{a}} \psi^* \left(\frac{t-b}{a} \right) dt \right|^2 \quad (5)$$

Wavelets can be classified as orthogonal, bi-orthogonal, semi-orthogonal or non-orthogonal wavelets.

The Morlet wavelet is a non-orthogonal wavelet given as:

$$\psi(t) = \frac{1}{\sqrt[4]{\pi}} \left(e^{j\omega_0 t} - e^{-\frac{\omega_0^2}{2}} \right) \cdot e^{-\frac{t^2}{2}}, \quad (6)$$

where ω_0 is the centre frequency.

In the case of $\omega_0 > 5$, Eq. (6) is simplified to (7):

$$\psi(t) = \frac{1}{\sqrt[4]{\pi}} e^{j\omega_0 t} \cdot e^{-\frac{t^2}{2}}. \quad (7)$$

Orthogonal wavelets can be constructed from polynomial spline functions or by solving for the filter coefficients such that its Fourier transform, frequency response function, satisfies orthogonality and moment conditions [14]. Such wavelets are the *Shannon*, *Meyer*, *Battle-Lemarie* and *Daubechies* compactly supported wavelets.

The Meyer wavelet is defined through the scaling function $\hat{\phi}(\omega)$ as:

$$\hat{\phi}(\omega) := \begin{cases} 1 & |\omega| < \frac{2\pi}{3} \\ \cos \left[\frac{\pi}{2} \nu \left(\frac{3}{2\pi} |\omega| - 1 \right) \right] & \frac{2\pi}{3} \leq |\omega| \leq \frac{4\pi}{3} \\ 0 & \text{otherwise} \end{cases}, \quad (8)$$

where $\nu(\omega)$ is a tapering function.

2.2 Matching Wavelet to a Given Signal

Chapa and *Rao* worked out a method that is capable of designing *Meyer wavelets* to match band-limited signals. Their method directly matches a wavelet to the signal. The method requires some conditions on the wavelet spectrum amplitude and phase. They define the theorem of band-limited scaling function and the band-limited wavelet as necessary and sufficient conditions for an Orthonormal Multiresolution Analysis (OMRA). They introduce a method by which a wavelet in OMRA is expressed in terms of scaling function. They also derive constraints on the structure of the wavelet phase. The matching algorithm is sub-optimal in the sense that it matches the wavelet amplitude and phase independently [11].

2.2.1 Matching Wavelet Amplitude

The starting point of amplitude matching algorithm is the discrete form of the *refinement equation* by which a scaling function can be expressed from a wavelet:

$$\left| \hat{\phi}\left(\frac{\pi k}{2^l}\right) \right|^2 = \sum_{p=0}^l \left| \hat{\psi}\left(\frac{2\pi k}{2^p}\right) \right|^2 \quad (9)$$

They define the condition on the wavelet spectra $\hat{\psi}(\omega)$ to guarantee orthonormality (Eq. (11)) and an error function (Eq. (10)) as:

$$E(Y, a) = \int_{\frac{2\pi}{3}}^{\frac{8\pi}{3}} [W(\omega) - aY(\omega)]^2 d\omega, \quad (10)$$

where:

$$W(\omega) = |F(\omega)|^2$$

$$Y(\omega) = |\hat{\psi}(\omega)|^2$$

a = scaling coefficient.

It is shown in [11] that this wavelet construction is Meyer's spectrum amplitude construction exactly. The algorithm searches for the extreme value of cost function in a discrete form. Using the symmetric property of the wavelet function, their design equation can be expressed in the form:

$$\sum_{p=0}^l \sum_{m=-\infty}^{+\infty} Y\left(\frac{2^l}{2^p}(k + 2^{l+1}m)\right) = 1, \quad (11)$$

where:

$$2^{l-1} / 3 < \left| \frac{2^l}{2^p} (k + 2^{l+1} m) \right| < 2^{l+2} / 3 \quad (12)$$

$\mathbf{1}$ = vector of all ones.

Equation (11) can be rewritten in matrix form as

$$\mathbf{A} \mathbf{Y} = \mathbf{1}, \quad (13)$$

where:

$$\mathbf{A} = \{ \alpha_{ij} \in \{0, 1, 2\}; \quad i = 1, \dots, L; \quad j = 1, \dots, 2^l \}$$

This amplitude matching algorithm is a constrained optimization problem that can be solved by *Lagrange multipliers* [11].

The error function (Eq. (10)) can be given by

$$E = \frac{(\mathbf{W} - a\mathbf{Y})^T (\mathbf{W} - a\mathbf{Y})}{\mathbf{W}^T \mathbf{W}}. \quad (14)$$

Matching amplitudes are given in the form of

$$\mathbf{Y} = \frac{1}{a} \mathbf{W} + \mathbf{A}^T (\mathbf{A} \mathbf{A}^T)^{-1} \left(\mathbf{I} - \frac{1}{a} \mathbf{A} \mathbf{W} \right), \quad (15)$$

where

$$a = \frac{\mathbf{I}^T (\mathbf{A} \mathbf{A}^T)^{-1} \mathbf{A} \mathbf{W}}{\mathbf{I}^T (\mathbf{A} \mathbf{A}^T)^{-1} \mathbf{I}}. \quad (16)$$

2.2.2 Matching Wavelet Phase

The phase matching algorithm is similar to the amplitude matching since it is based on MSE criteria, but instead of signal phase, it uses group delay. The group delay of a signal is defined as the first order, negative derivative of the phase

$$\tau(\omega) = - \frac{d\theta(\omega)}{d\omega}.$$

There are specific constraints on the structure of the wavelet phase (Eq. (20)), which is expressed in terms of the phase of the scaling function [11]:

$$\theta_{\Psi}(\omega) = -\frac{\omega}{2} - \theta_{\Phi}(\omega + 2\pi) + \theta_{\Phi}\left(\frac{\omega}{2} + \pi\right) + \theta_{\Phi}\left(\frac{\omega}{2}\right), \quad (17)$$

where θ_{Φ} is the phase of scaling function $\phi(\omega)$ and θ_{Ψ} is the phase of wavelet $\Psi(\omega)$.

The wavelet phase is a symmetric, 2π periodic, even function. The method models one period of the negative of the group delay, denoted by $\lambda_T(\omega)$ as a polynomial of order R [11]:

$$\lambda_T(\omega) = \sum_{r=0}^{\frac{R}{2}} c_r \omega^{2r} \Pi\left(\frac{\omega}{2}\right), \quad (18)$$

where

$$\lambda(\omega) = -\tau(\omega)$$

$$\Pi(\omega) = \begin{cases} 1, & -\frac{1}{2} \leq \omega < \frac{1}{2} \\ 0, & \text{otherwise} \end{cases}$$

c_r = polynomial coefficients.

By replicating one period of the group delay at every 2π interval, the group delay of the wavelet is modelled as the 2π periodic polynomial of order R [11]:

$$\lambda(\omega) = \sum_{k=-\infty}^{+\infty} \lambda_T(\omega - 2\pi k) = \sum_{k=-\infty}^{+\infty} \sum_{r=0}^{R/2} c_r (\omega - 2\pi k)^{2r} \Pi\left(\frac{\omega - 2\pi k}{2\pi}\right) \quad (19)$$

The discrete form of Eq. (19) can be written as (20), where $\Delta\omega = \frac{2\pi}{T}$, $P = \frac{N}{T}$,

N is the number of samples, $-\frac{N}{2} \leq n < \frac{N}{2}$.

$$\lambda(n) = \sum_{r=0}^{R/2} c_r \sum_{k=-P/2}^{P/2-1} (n - kT)^{2r} \Pi\left(\frac{n - kT}{T}\right) \quad (20)$$

Using vector notation, the group delay is expressed as

$$\lambda = \mathbf{B} \mathbf{c}, \quad (21)$$

where

$$\mathbf{b}_{n,r} = \sum_{k=-\frac{P}{2}}^{\frac{P}{2}} (n - kT)^{2r} \Pi\left(\frac{n - kT}{T}\right). \quad (22)$$

Negatives of the group delays Λ_ψ and Λ_ϕ can be expressed in terms of $\lambda(\omega)$.

Applying $\Gamma_\psi(\omega) = \Lambda_\psi + \frac{1}{2}$ substitution [11], we get

$$\Gamma_\psi = \mathbf{D}_\psi c = -\frac{1}{2} \mathbf{B}_{\left(\frac{q+T}{2}\right)} + \sum_{m=2}^{\infty} 2^{-m} \mathbf{B}_{\frac{q}{2^m}}, \quad (23)$$

$$\Gamma_\phi = \mathbf{D}_\phi c = \sum_{m=1}^{\infty} 2^{-m} \mathbf{B}_{\frac{q}{2^m}}. \quad (24)$$

The matching process minimizes the weighted error (Eq. (25)) between the group delay of the wavelet Γ_ψ and the desired signal Γ_F [11]. The approximation is performed only in the pass-band [11], thus a weighting function $\Omega(n)$ is calculated from the result of the amplitude matching process:

$$\gamma = \sum_{n=-N/2}^{N/2-1} \left(\Omega(n) (\Gamma_F(n) - \Gamma_\psi(n)) \right)^2, \quad (25)$$

where

$$\Omega(n) = \frac{Y(n)}{\sum Y(n)}.$$

The optimal values of the polynomial coefficients can be obtained by solving

$$\nabla_c \gamma = 0. \quad (26)$$

2.2.3 Matching a Wavelet to Signal Model of Bearing Vibration

We used the above-mentioned algorithm to create a matching wavelet to the transient vibration generated by an inner race fault – a pitting or spalling formulation – in a deep groove ball bearing. It is shown in [16] that the rolling elements generate a series of amplitude modulated transient pulses when they pass over the fault. The amplitude modulation is caused by load distribution; that is, the closer the fault is located to the load zone, the higher the amplitude of the transient is. One of these impulses can be described as:

$$y(t) = A \cdot t^n \cdot e^{-C \cdot t} \cdot \sin(\omega_0 \cdot t), \quad t \in [0, \infty) \quad \omega_n = 2\pi f_n \quad (27)$$

where f_n is the n^{th} natural frequency of bearing system, C is a damping factor, A is the initial amplitude and n is an exponent influencing the rise time of the transient.

In order to create a new wavelet basis function, we used 512 samples of time domain data of a transient vibration signal (Fig. 1) described by Eq. (27) with $A=68.74$, $n=1.851$, $C=6.78$, and $\omega_0=18.85$.

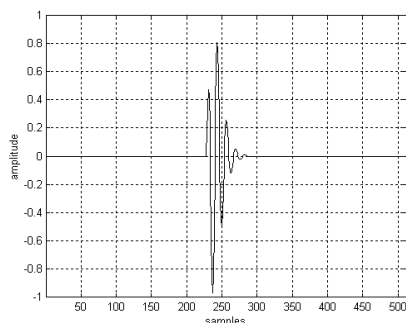


Figure 1

Transient signal model ($A=68.74$, $n=1.851$, $C=6.78$, $\omega_0=18.85$)

We applied *Chapa and Rao's* amplitude and phase matching algorithm [11] on the transient signal model of the bearing vibration. The results are shown in Fig. 2. The new wavelet amplitude spectra match the amplitude spectra of the transient very well in the passband. The MSE of the matching is 0.011. For phase matching we used a 16th-degree polynomial. The matched group delay shows satisfactory characteristics in the passband. These plots are very similar to the plots introduced in [11], since the applied transient pulses are similar in both cases.

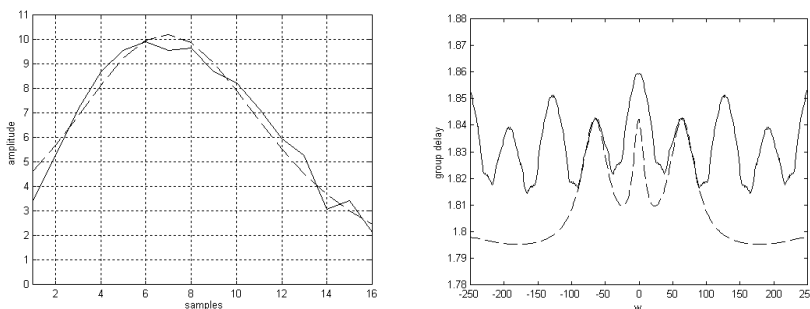


Figure 2

Matching the wavelet amplitude and the group delay in the neighbourhood of the passband (matched wavelet amplitude spectra and group delay: continuous line; amplitude spectra and group delay of the transient: dashed line)

Combining the amplitude and phase spectra together, it is possible to obtain the time domain form of wavelet function [11]. It can be seen in Fig. 3 that the new wavelet adequately fits the original transient data.

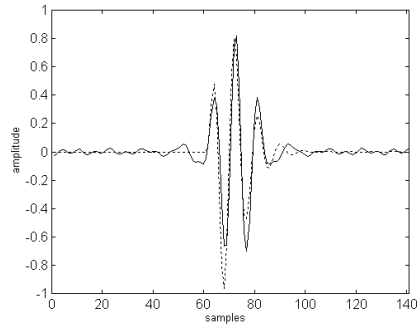


Figure 3

Transient signal and the new wavelet in the time domain (transient signal: dotted line; wavelet: continuous line)

Since the new wavelet basis cannot be given in closed form we gave the filter coefficients in Table 1. The time- and frequency behaviour of the wavelet and the impulse responses of filters are shown in Figs. 4 and 5.

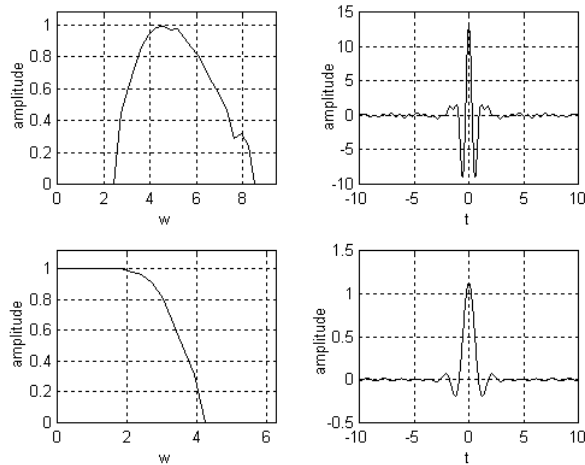


Figure 4

The new wavelet and scaling function in the frequency and time-domain

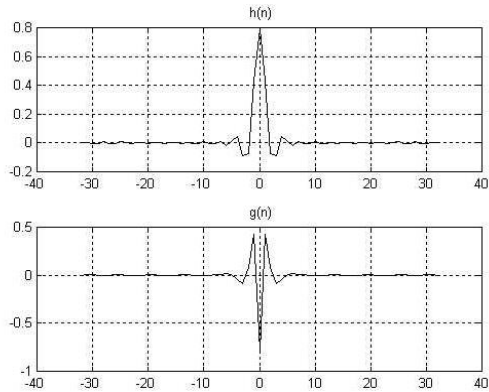


Figure 5

The impulse responses of filters corresponding to new wavelet

From Table 1, one can conclude that these *Quadrature Mirror Filters* (QMF) have no compact support. Their effective support is approximately $[-10, 10]$. But by limiting the filter coefficients to this range, we cannot reconstruct the original data from the wavelet coefficients completely. The fewer coefficients we use, the more error we get during reconstruction. The opposite is also true; that is, more coefficients are needed to reconstruct the original data from the wavelet coefficients. Since this wavelet was created to match the vibration signal of a bearing with a specified fault, it can be used to detect this kind of defect in a bearing.

Table 1

The new wavelet $\Psi(t)$ filter coefficients $\{h[k]\}$ and $\{g[k]\}$

k	h(n)	g(n)
± 0	0.7948	-0.7948
± 1	0.4260	0.4260
± 2	-0.0760	0.0760
± 3	-0.0872	-0.0872
± 4	0.0474	-0.0474
± 5	0.0115	0.0115
± 6	-0.0166	0.0166
± 7	0.0068	0.0068
± 8	-0.0028	0.0028
± 9	-0.0024	-0.0024
± 10	0.0069	-0.0069
± 11	-0.0045	-0.0045
± 12	-0.0012	0.0012

3 Application of Matched and Morlet Wavelet for Detection of Spalling in a Bearing

We examined the applicability of the matched and the Morlet wavelet for detecting pitting formulation on the inner race of a deep-groove ball bearing. Our aim was to detect transient pulses generated in the bearing to indicate the presence of a pitted raceway, i.e. the time of the last possible maintenance before catastrophic failure.



Figure 6

A single row deep-groove ball-bearing of type 6209 used as test specimen



Figure 7

Pitting formulated on the inner raceway during endurance test

We used two 6209-type, single-row, deep-groove, radial ball bearings as test specimens (Fig. 6). One of them was earlier subjected to an endurance test. As a consequence of the repetitive load on the bearing elements, pitting formed on the inner ring of the bearing (Fig. 7). The other bearing was free of faults.

For the test rig we used a turning machine of E1N type. The inspected bearing was mounted on a shaft which was fixed in the chuck. We used a rod fixed in the tool post as the support. The rotating nature of the tool post made it possible to apply radial load on the outer ring of bearing, where the force was set to be perpendicular to the rotating shaft. Our primary goal was to minimize the force/vibration transmission path, since noise can come from a number of different sources. Mechanical noises can be eliminated by properly set up measurement device configuration, while electrical noises usually come from the test rig. A portion of the outer ring of bearing was machined by grinding. An accelerometer of KISTLER 8702B50 type was attached to its flat area by beeswax.

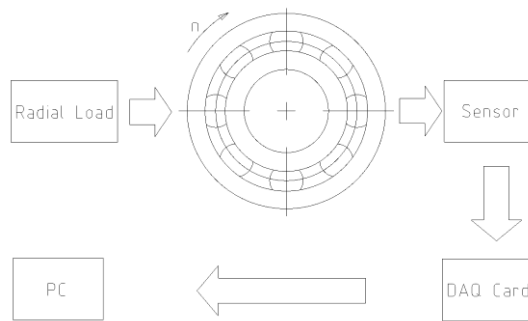


Figure 8
The measurement setup

For data acquisition (Fig. 8) we used the following devices:

- HAMEG, HM507 analog-digital oscilloscope, 100 Ms/s real-time sampling rate,
- KISTLER accelerometer 8702B50,
- KISTLER 5108 charge amplifier,
- PCI 6063E PCMCIA DAQ card, 500 ks/s sampling rate.

The DAQ card was controlled by software developed under the NI LabWindows/CVI programming environment. Validation of our software was performed using a HITACHI VG-4429 function generator and digital oscilloscope.

Sampling was performed at a constant inner ring speed of 1812 min^{-1} . This value satisfies the specifications of American ANSI [17] and German DIN [18] standards ($1800 \text{ min}^{-1} \pm 2\%$) concerning bearing vibration measurements. The outer ring was stationary, as it delivered radial load. The sampling frequency was set to be 20 kHz since the accelerometer's frequency range ends at 10 kHz. The gain was set to unity.

The time series data of sampled bearing vibration signal is shown in Fig. 9 and Fig. 10.

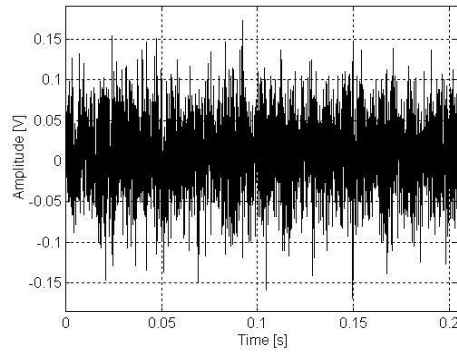


Figure 9

Vibration data of the good bearing, $f_{\text{sampling}}=20$ kHz

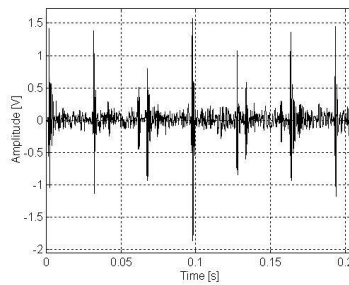


Figure 10

Vibration data of the pitted bearing, $f_{\text{sampling}}=20$ kHz

The vibration data are the results of two distinct measurements, where we could not provide exactly the same load on the bearings. Therefore one cannot make a final decision comparing the numerical values of vibration amplitudes. These figures qualitatively indicate the vibration signals. However, the difference is obviously shown. One can notice the repetitive transient pulses with an average periodicity of 30 ms. This variation in repetition time is the result of the motion of the bearing elements, which are rolling and sliding.

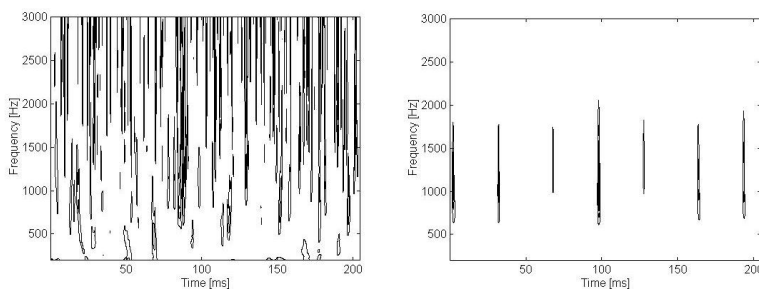


Figure 11

Time-frequency distribution of signal energy of good (left) and pitted bearing (right) using the matched wavelet

The scalogram of the good bearing clearly shows the time-frequency location of transients. These pulses do not appear at each rotation. They seem to be random signals that might come from the test rig. In contrast, periodically repeating transient pulses are clearly seen on the scalogram of the pitted bearing. Their time-period and frequency can be numerically given. The application of this method reduces the incorrect evaluation of vibration data and can be a valuable supplement to conventional condition monitoring methods.

We calculated the scalogram of the same vibration data using the Morlet wavelet as well (Fig. 12). The choice of this wavelet is obvious, since many authors report its successful application to bearing vibration analysis [5, 7, 8, 15, 19-21].

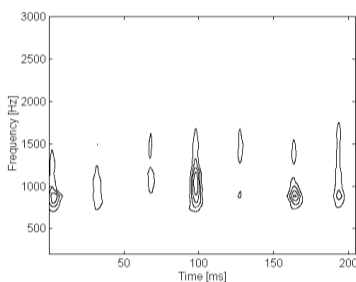


Figure 12

Time-frequency distribution of signal energy of the pitted bearing using the Morlet wavelet

Comparing the two time-frequency representations that were calculated using the matched wavelet (Fig. 11) and the Morlet wavelet (Fig. 12), we notice that the Morlet wavelet provided a more realistic result. The Morlet wavelet gave better energy concentration. The scalogram calculated by the matching wavelet provided better time localization, but its frequency localization is less accurate than that of the Morlet wavelet. This raises the question of the cause of better representation using the Morlet wavelet, since the matched wavelet was designed using the signal model of this type of failure in bearings.

4 Examination of the Wavelets and the Signal Model

The time and frequency domain plots of the matched and the Morlet wavelet are shown in Figs. 13 and 14. The wavelet amplitudes were normalized to facilitate comparison.

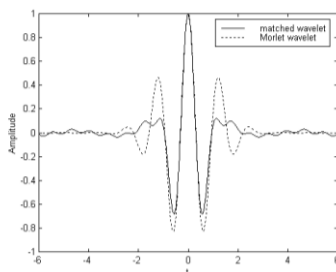


Figure 13

Time plot of the matched and the Morlet wavelet

The matched and the Morlet wavelet follow almost the same pattern between the $[-1.5, 1.5]$ time interval, where the matched wavelet approximates the Morlet wavelet. Most of the signal energy is concentrated in this area.

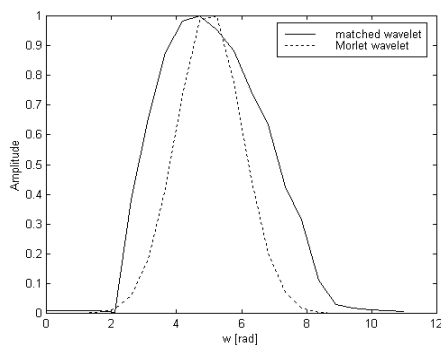


Figure 14

Amplitude spectrum of the matched and Morlet wavelet

The frequency domain form of these wavelets shows their bandpass behaviour. The bandwidth of the Morlet wavelet is narrower than that of the matched; that is, the Morlet wavelet concentrates more of the signal energy around the centre frequency. The wider bandwidth gives shorter time extent; thus the time localization ability of the matched wavelet is better. This is also clearly seen in Fig. 11. and Fig. 12.

Comparing the real part of the complex Morlet wavelet in Eq. (7) to the signal model of transient vibration produced by a pitting formulation on the inner race of

a deep groove ball bearing (Eq. (27)), one can notice the similarity. These waveforms are shown in Figs. 15 and 16. The resemblance is clearly seen.

The frequency domain form of these waveforms shows their bandpass behaviour (Fig. 16). The phase characteristics of these waveforms are also similar except for the dilation caused by unwrapping their phase.

Several papers report the successful application of the Morlet wavelet in the field of bearing vibration analysis without clarifying the reason for their choice. Since the signal model of transient pulses generated by pitting on the inner raceway of a deep-groove ball bearing was very similar to the Morlet wavelet, and since the wavelet transform and its derived energy distributions give more output when the signal and the analysing wavelet are similar, we can conclude that for analysing exponentially or near-exponentially damped vibration responses like the bearing vibration signal caused by pitting formulation on the inner raceway of a deep-groove ball bearing, the Morlet wavelet is a reasonable choice that is sure to yield good results.

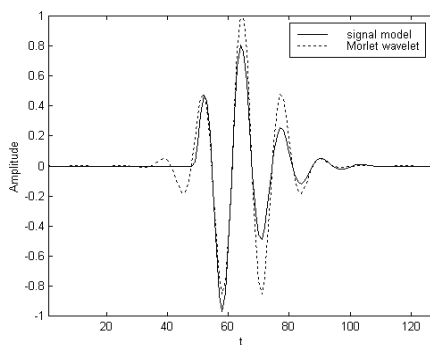


Figure 15

Signal model of bearing vibration (Eq. (30)) and the Morlet wavelet in the time domain

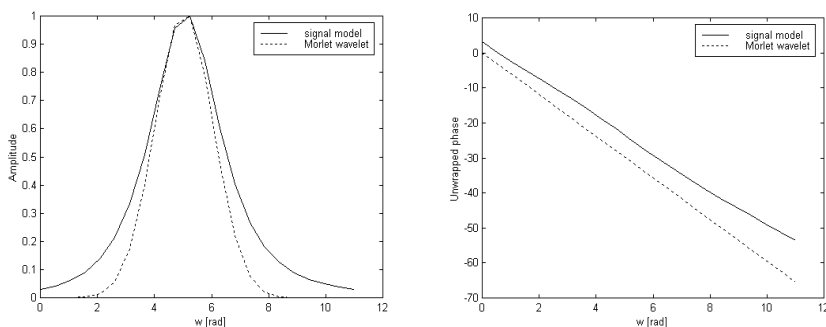


Figure 16

Amplitude and phase spectrum of the signal model and Morlet wavelet

Conclusions

This paper showed the creation of a new wavelet that matches the transient vibration response generated by pitting on the inner race of a deep-groove ball bearing. Wavelet creation is based on *Chapa and Rao's* method, where the Meyer wavelet amplitude and phase spectra are matched independently to the signal. It was shown that the new wavelet can be used for detecting transient pulses generated in a bearing. We compared the results with those calculated using the Morlet wavelet since many application reported its successful application. We found the Morlet wavelet superior to this matched wavelet in representing transient signals of bearing vibration where pitting is formulated in the inner raceway of a deep groove ball bearing. It was shown that the signal model of this kind of bearing failure is very similar to the Morlet wavelet; thus its scalogram gives a more accurate time-frequency representation than the Meyer wavelet.

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